

LES of Turbulent Flows: Lecture 6

(ME EN 7960-003)

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Homework #2

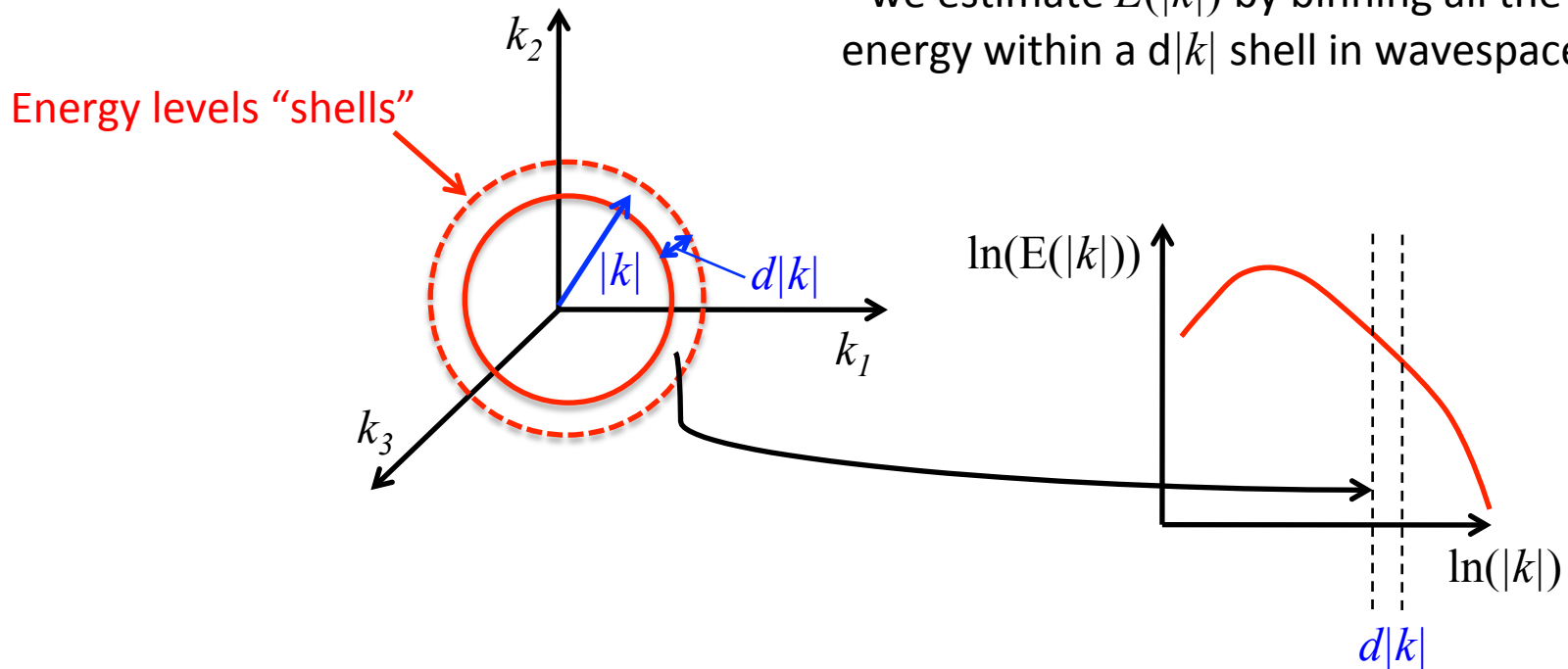
- Goal of homework assignment is to calculate 3D energy spectrum from isotropic data and to perform 3D filtering on a 3D isotropic turbulence dataset

- **Calculating the 3D energy spectrum**

- plot: $E(|k|)$ vs. $|k| = [k_1^2 + k_2^2 + k_3^2]^{1/2}$

- $E(|k|) = |\hat{u}_{\vec{k}}|^2 + |\hat{v}_{\vec{k}}|^2 + |\hat{w}_{\vec{k}}|^2$

- we estimate $E(|k|)$ by binning all the energy within a $d|k|$ shell in wavespace



LES filtered Equations for incompressible flow

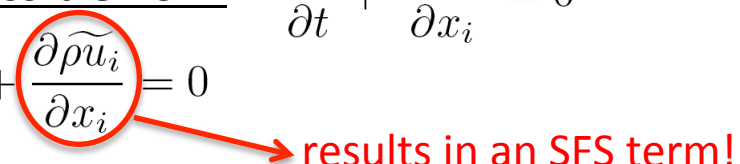
- Mass: $\frac{\partial \tilde{u}_i}{\partial x_i} = 0$
- Momentum: $\frac{\partial \tilde{u}_i}{\partial t} + \frac{\partial \tilde{u}_i \tilde{u}_j}{\partial x_j} = -\frac{\partial \tilde{p}}{\partial x_i} + \frac{1}{Re} \frac{\partial^2 \tilde{u}_i}{\partial x_j^2} - \frac{\partial \tau_{ij}}{\partial x_j} + F_i$
- Scalar: $\frac{\partial \tilde{\theta}}{\partial t} + \frac{\partial \tilde{u}_i \tilde{\theta}}{\partial x_i} = \frac{1}{Sc Re} \frac{\partial^2 \tilde{\theta}}{\partial x_i^2} - \frac{\partial q_i}{\partial x_i} + Q$
- Resolved Energy: $\frac{\partial \tilde{E}_f}{\partial t} + \tilde{u}_i \frac{\partial \tilde{E}_f}{\partial x_i} = -\frac{1}{\rho} \frac{\partial \tilde{u}_i \tilde{P}}{\partial x_i} - \frac{\partial \tilde{u}_i \tau_{ij}}{\partial x_j} - 2\nu \frac{\partial \tilde{u}_i \tilde{S}_{ij}}{\partial x_j} - \varepsilon_f - \Pi$
- SFS stress: $\tau_{ij} = \widetilde{u_i u_j} - \tilde{u}_i \tilde{u}_j$
- SFS flux: $q_j = \widetilde{u_j \theta} - \tilde{u}_j \tilde{\theta}$
- SFS Dissipation: $\Pi = -\tau_{ij} \tilde{S}_{ij}$

Filtering the compressible N-S equations

- What happens when we apply a filter to the compressible N-S equations?

-Conservation of Mass for compressible flow: $\frac{\partial \rho}{\partial t} + \frac{\partial \rho u_i}{\partial x_i} = 0$

• filtering each term => $\frac{\partial \tilde{\rho}}{\partial t} + \frac{\partial \tilde{\rho} \tilde{u}_i}{\partial x_i} = 0$

 results in an SFS term!

- How can we avoid having a SFS conservation of mass?

-Density weighted filtering:

- Formalized for compressible flow by Favre (Phys. Fluids, 1983) for ensemble statistics, a Favre (or density weighted) filter is defined by:

$$\bar{\phi} = \frac{\tilde{\rho} \tilde{\phi}}{\tilde{\rho}} \Rightarrow \tilde{\rho} \bar{\phi} = \tilde{\rho} \tilde{\phi}$$

where we note that as compressibility becomes less important $\bar{\phi} \rightarrow \tilde{\phi}$
and we can show that the conservation of mass becomes:

$$\frac{\partial \tilde{\rho}}{\partial t} + \frac{\partial \tilde{\rho} \tilde{u}_i}{\partial x_i} = 0$$

Filtering the compressible N-S equations

- We can use this to write the Favre filtered equations of motion (see Geurts pg 32-35 or Vreman et al. Applied Sci. Res. 1995 for details)

-Conservation of Mass:
$$\frac{\partial \tilde{\rho}}{\partial t} + \frac{\partial \tilde{\rho} \bar{u}_i}{\partial x_i} = 0$$

-Conservation of Momentum:

$$\frac{\partial \tilde{\rho} \bar{u}_i}{\partial t} + \frac{\partial \tilde{\rho} \bar{u}_i \bar{u}_j}{\partial x_j} + \frac{\partial \tilde{p}}{\partial x_i} - \frac{\partial \bar{\sigma}_{ij}}{\partial x_j} = -\frac{\partial \tilde{\rho} \tau_{ij}}{\partial x_j} + \frac{\partial}{\partial x_j} (\tilde{\sigma}_{ij} - \bar{\sigma}_{ij})$$

where the SFS terms are collected on the RHS of the equation and we now have both a resolved ($\bar{\sigma}_{ij}$) and SFS ($\tilde{\sigma}_{ij} - \bar{\sigma}_{ij}$) viscous contribution because $\mu = \mu(\bar{T})$ is a function of the Favre filtered temperature and

$$\tilde{\sigma}_{ij} = \left(\frac{2}{Re}\right) \mu(T) \left(S_{ij} - \frac{1}{3} \delta_{ij} \frac{\partial u_k}{\partial x_k} \right) \Rightarrow \text{nonlinear viscous stress tensor}$$

Recall: strain rate tensor

$$\bar{\sigma}_{ij} = \left(\frac{2}{Re}\right) \mu(\bar{T}) \left(\bar{S}_{ij} - \frac{1}{3} \delta_{ij} \frac{\partial \bar{u}_k}{\partial x_k} \right) \Rightarrow \text{“smooth” viscous stress tensor}$$

$$S_{ij} = \frac{1}{2} \left(\frac{\partial u_i}{\partial x_j} + \frac{\partial u_j}{\partial x_i} \right)$$

- The SFS stress tensor for the Favre filtered equations is given by

$$\tau_{ij} = \overline{u_i u_j} - \bar{u}_i \bar{u}_j$$

which is obtained from the nonlinear term $\Rightarrow \widetilde{\rho u_i u_j} = \tilde{\rho} \overline{u_i u_j} = \tilde{\rho} (\overline{u_i u_j} - \bar{u}_i \bar{u}_j)$

Filtering the compressible N-S equations

-Conservation of total energy: $\bar{e} \equiv$ Favre filtered total energy density $= \frac{\tilde{p}}{\gamma - 1} + \frac{1}{2} \tilde{\rho} \bar{u}_i \bar{u}_i$

$$\frac{\partial \bar{e}}{\partial t} + \frac{\partial}{\partial x_j} ((\bar{e} + \tilde{p}) \bar{u}_j) - \frac{\partial \bar{u}_i \bar{\sigma}_{ij}}{\partial x_j} + \frac{\partial \bar{q}_j}{\partial x_j} = -a_1 - a_2 - a_3 + a_4 + a_5 - a_6$$

where the LHS contains the SFS terms created using the procedure used in Lecture 5 for τ_{ij}

$a_1 = \bar{u}_i \frac{\partial \tilde{p} \tau_{ij}}{\partial x_j}$	→ kinetic energy transferred from resolved to SFSs
$a_2 = \frac{1}{\gamma - 1} \frac{\partial (\widetilde{p u_j} - \tilde{p} \bar{u}_j)}{\partial x_j}$	→ pressure velocity SFS term (effect of SFS turbulence on the conduction of heat at resolved scales)
$a_3 = p \frac{\partial \widetilde{u_j}}{\partial x_j} - \tilde{p} \frac{\partial \bar{u}_j}{\partial x_j}$	→ compressibility effects (vanishes for incompressible)
$a_4 = \widetilde{\sigma_{ij} \frac{\partial u_i}{\partial x_j}} - \tilde{\sigma}_{ij} \frac{\partial \bar{u}_i}{\partial x_j}$	→ conversion of SFS kinetic energy to internal energy by viscous dissipation
$a_5 = \frac{\partial (\tilde{\sigma}_{ij} \bar{u}_i - \bar{\sigma}_{ij} \bar{u}_i)}{\partial x_j}$	→ SFS viscous stress term
$a_6 = \frac{\partial (\tilde{q}_j - \bar{q}_j)}{\partial x_j}$	→ SFS heat flux term (Note q_j is the heat flux vector)

Typically assumptions that $\tilde{\sigma}_{ij} - \bar{\sigma}_{ij} \approx 0$ and $\tilde{q}_j - \bar{q}_j \approx 0$ are made eliminating a_5 and a_6 .